# A Comment on the Vibrational Analysis for $\mathrm{C}_{60}$ and Other Fullerenes 

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Recently Nakamoto and McKinney (1) provided a symmetry analysis of the vibrational modes of $\mathrm{C}_{60}$ and other fullerenes. I would like to supplement their presentation with another approach, in widespread use, which records the number of atoms that are unmoved by each symmetry operation, yielding the reducible representation $\Gamma_{u m a}$ (2). This is particularly easy to do for $\mathrm{C}_{60}$ because only the identity operation and the 15 symmetry planes leave atoms unmoved, 60 and 4 , respectively, as is shown in the accompanying table.

Multiplication of $\Gamma_{\text {uma }}$ by the representation for translation in the x-, y-, and z-directions ( $\mathrm{T}_{1 \mathrm{u}}$ in $\mathrm{I}_{\mathrm{h}}$ symmetry) yields, $\Gamma_{\text {tot }}$, the reducible representation for the 180 degrees of freedom of the $\mathrm{C}_{60}$ molecule. $\Gamma_{\text {tot }}$ is then decomposed into the equivalent linear combination of $I_{h}$ irreducible representations by the usual method as shown in the table. The symmetry of the vibrational modes is found by subtracting $\Gamma_{\text {trans }}$ and $\Gamma_{\text {rot }}$ from $\Gamma_{\text {tot }}$.

$$
\Gamma_{\text {vib }}=2 \mathrm{~A}_{\mathrm{g}}+3 \mathrm{~T}_{1 \mathrm{~g}}+4 \mathrm{~T}_{2 \mathrm{~g}}+6 \mathrm{G}_{\mathrm{g}}+8 \mathrm{H}_{\mathrm{g}}+\mathrm{A}_{\mathrm{u}}+4 \mathrm{~T}_{1 \mathrm{u}}+5 \mathrm{~T}_{2 \mathrm{u}}+6 \mathrm{G}_{\mathrm{u}}+7 \mathrm{H}_{\mathrm{u}}
$$

IR active modes have the same symmetry as $\mathrm{x}, \mathrm{y}$, and z . Raman active modes have the symmetry of the quadratic forms, $\mathrm{x}^{2}$, xy , etc. Thus, we find four IR active modes and ten Raman active modes in agreement with experimental spectroscopic evidence. (1) The symmetry of the stretching modes can be determined by examining the behavior of the carbon-carbon bonds, $\Gamma_{\text {bonds }}$, under the symmetry operations of the $I_{h}$ group. In addition, the symmetry of the $\pi$-electron density can be easily studied because $\Gamma_{\pi}=\Gamma_{\text {uma }}$. The combination of knowing the irreducible representations contributing to $\Gamma_{\text {vib }}$ and $\Gamma_{\pi}$ allows one to do an in depth analysis of electronic spectrum of $\mathrm{C}_{60}$ utilizing the mechanism of vibronic coupling (3).

It has been shown that this method of recording the number of unmoved atoms is ideally
suited for computer programming environments that have matrix and vector algebra capability (4). It is easy, for example, to prepare a Mathcad worksheet for any finite point group, which can then serve as a template for any molecule with that symmetry. The only thing that changes from one molecule to another within the same point group is the vector representing $\Gamma_{\text {uma }}$.

## Literature Cited:

1. Nakamoto, K.; McKinney, M. A. J. Chem. Educ. 2000, 77, 775.
2. Harris, D. C.; Bertolucci, M. D. Symmetry and Spectroscopy: An Introduction to

Vibrational and Electronic Spectroscopy ; Dover Publications, Inc.: New York, 1989, p. 141.
3. Leach, S.; et al. Chem. Phys.1992, 160, 451.
4. Rioux, F. J. Chem. Educ: Soft. 1998, Issue 9801MW.

| Character Table for the Icosahedral Point Group |  |  |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{I}_{\mathrm{h}}$ | E | $12 \mathrm{C}_{5}$ | $12 \mathrm{C}_{5}{ }^{2}$ | $20 \mathrm{C}_{3}$ | $15 \mathrm{C}_{2}$ | i | $12 \mathrm{~S}_{10}$ | $12 \mathrm{~S}_{10}{ }^{2}$ | $20 \mathrm{~S}_{6}$ | 150 | $\mathrm{h}=120$ |
| $\mathrm{A}_{\mathrm{g}}$ | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | $\mathrm{x}^{2}+\mathrm{y}^{2}+\mathrm{z}^{2}$ |
| $\mathrm{T}_{1 \mathrm{~g}}$ | 3 | 1.618 | -0.618 | 0 | -1 | 3 | 1.618 | -0.618 | 0 | -1 | $\mathrm{R}_{\mathrm{x}}, \mathrm{R}_{\mathrm{v}}, \mathrm{R}_{\mathrm{z}}$ |
| $\mathrm{T}_{2 \mathrm{~g}}$ | 3 | -0.618 | 1.618 | 0 | -1 | 3 | -0.618 | 1.618 | 0 | -1 |  |
| $\mathrm{G}_{\underline{g}}$ | 4 | -1 | -1 | 1 | 0 | 4 | -1 | -1 | 1 | 0 |  |
| $\mathrm{Hg}_{\underline{g}}$ | 5 | 0 | 0 | -1 | 1 | 5 | 0 | 0 | -1 | 1 | xy,xz,yz |
| $\mathrm{A}_{\mathrm{u}}$ | 1 | 1 | 1 | 1 | 1 | -1 | -1 | -1 | -1 | -1 |  |
| $\mathrm{T}_{1 \mathrm{u}}$ | 3 | 1.618 | -0.618 | 0 | -1 | -3 | -1.618 | 0.618 | 0 | 1 | $\mathrm{x}, \mathrm{y}, \mathrm{z}$ |
| $\mathrm{T}_{2 \mathrm{u}}$ | 3 | -0.618 | 1.618 | 0 | -1 | -3 | 0.618 | -1.618 | 0 | 1 |  |
| $\mathrm{G}_{\mathrm{u}}$ | 4 | -1 | -1 | 1 | 0 | -4 | 1 | 1 | -1 | 0 |  |
| $\mathrm{H}_{\mathrm{u}}$ | 5 | 0 | 0 | -1 | 1 | -5 | 0 | 0 | 1 | -1 |  |
| $\Gamma_{\text {uma }}$ | 60 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 4 |  |
| $\Gamma_{\text {tot }}=\Gamma_{\text {uma }} \cdot \mathrm{T}_{\text {lu }}$ | 180 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 4 |  |
| $\Gamma_{\text {tot }} \cdot \mathrm{Ag}_{\mathrm{g}}$ | 180 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 60 | $\div \mathrm{h}=2$ |
| $\Gamma_{\text {tot }} \mathrm{T}_{1 \mathrm{~g}}$ | 540 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | -60 | $\div \mathrm{h}=4$ |
| $\Gamma_{\text {tot }} \cdot \mathrm{T}_{2 \mathrm{~g}}$ | 540 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | -60 | $\div \mathrm{h}=4$ |
| $\Gamma_{\text {tot }} \cdot \mathrm{G}_{\mathrm{g}}$ | 720 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | $\div \mathrm{h}=6$ |
| $\Gamma_{\text {tot }} \cdot \mathrm{Hg}_{\mathrm{g}}$ | 900 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 60 | $\div \mathrm{h}=8$ |
| $\Gamma_{\text {tot }} \cdot \mathrm{A}_{\mathrm{u}}$ | 180 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | -60 | $\div \mathrm{h}=1$ |
| $\Gamma_{\text {tot }} \mathrm{T}_{\text {lu }}$ | 540 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 60 | $\div \mathrm{h}=5$ |
| $\Gamma_{\text {tot }} \mathrm{T}_{2 \mathrm{u}}$ | 540 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 60 | $\div \mathrm{h}=5$ |
| $\Gamma_{\text {tot }} \cdot \mathrm{G}_{\mathrm{u}}$ | 720 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | $\div \mathrm{h}=6$ |
| $\Gamma_{\text {tot }} \cdot \mathrm{H}_{\mathrm{u}}$ | 900 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | -60 | $\div \mathrm{h}=7$ |

