

Another GHZ Example Using Spin-1/2 Particles

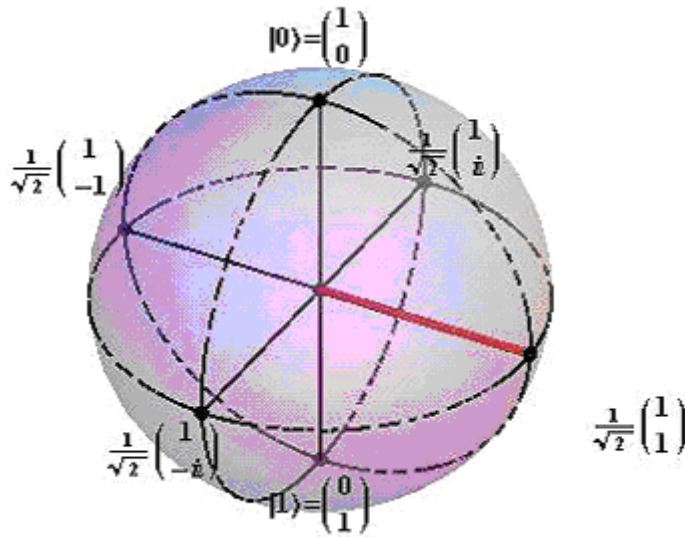
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This exercise explores the outcomes of measurements on the three spin-1/2 entangled state, Ψ_{ABC} , highlighted below. It represents a lean GHZ protocol developed by H. J. Bernstein. Because much has been published on the GHZ protocol, I'm going to get right to the point without a lot of commentary.

The eigenstates for spin-up and spin-down in the z-, x- and y-directions in vector format are as follows.

$$Z_u := \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad Z_d := \begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad X_u := \frac{1}{\sqrt{2}} \cdot \begin{pmatrix} 1 \\ 1 \end{pmatrix} \quad X_d := \frac{1}{\sqrt{2}} \cdot \begin{pmatrix} 1 \\ -1 \end{pmatrix} \quad Y_u := \frac{1}{\sqrt{2}} \cdot \begin{pmatrix} 1 \\ i \end{pmatrix} \quad Y_d := \frac{1}{\sqrt{2}} \cdot \begin{pmatrix} 1 \\ -i \end{pmatrix}$$

These spin states are also shown on a Bloch sphere.



Tensor multiplication using Mathcad commands will be used to form the initial entangled spin state and subsequent measurement states. These results can be easily verified by hand calculation.

$$\Psi(a, b, c) := \text{submatrix} \left[\text{kroncker} \left[\text{augment} \left[a, \begin{pmatrix} 0 \\ 0 \end{pmatrix} \right], \text{kroncker} \left[\text{augment} \left[b, \begin{pmatrix} 0 \\ 0 \end{pmatrix} \right], \text{augment} \left[c, \begin{pmatrix} 0 \\ 0 \end{pmatrix} \right] \right] \right], 1, 8, 1, 1 \right]$$

Initial State: $\Psi_{ABC} := \frac{1}{\sqrt{2}} \cdot (\Psi(Y_u, Y_u, Y_u) - \Psi(Y_d, Y_d, Y_d))$ $\Psi_{ABC}^T = (0 \ 0.5i \ 0.5i \ 0 \ 0.5i \ 0 \ 0 \ -0.5i)$

Given this initial state the probabilities for the various measurement outcomes for spin in the z-direction are calculated.

$$\left(\left| \Psi(Z_u, Z_u, Z_u)^T \cdot \Psi_{ABC} \right| \right)^2 = 0 \quad \left(\left| \Psi(Z_u, Z_u, Z_d)^T \cdot \Psi_{ABC} \right| \right)^2 = 0.25 \quad \left(\left| \Psi(Z_u, Z_d, Z_u)^T \cdot \Psi_{ABC} \right| \right)^2 = 0.25$$

$$\left(\left| \Psi(Z_d, Z_u, Z_u)^T \cdot \Psi_{ABC} \right| \right)^2 = 0.25 \quad \left(\left| \Psi(Z_u, Z_d, Z_d)^T \cdot \Psi_{ABC} \right| \right)^2 = 0 \quad \left(\left| \Psi(Z_d, Z_u, Z_d)^T \cdot \Psi_{ABC} \right| \right)^2 = 0$$

$$\left(\left| \Psi(Z_d, Z_d, Z_u)^T \cdot \Psi_{ABC} \right| \right)^2 = 0 \quad \left(\left| \Psi(Z_d, Z_d, Z_d)^T \cdot \Psi_{ABC} \right| \right)^2 = 0.25$$

These calculations predict that the following spin states are observed: $Z_u Z_u Z_d$, $Z_u Z_d Z_u$, $Z_d Z_d Z_u$ and $Z_d Z_d Z_d$. It is easily seen that the measurement of any two spin-states allows the prediction of the third without the need for a measurement. If two spins have the same value (uu or dd) then the third is d. If two spin states have different values (ud or du) then the third is u. Einstein, Podolsky and Rosen (EPR) wrote in their famous 1935 *Physical Review* paper,

"If, without in any way disturbing a system, we can predict with certainty (i.e. with probability equal to unity) the value of a physical quantity, then there exists an element of reality corresponding to this physical quantity."

According to this definition spin in the z-direction is an "element of reality." In other words, it has a definite value independent of measurement.

Next we calculate the quantum mechanical predictions for the measurement of spin in the x-direction on two of the spins and in the z-direction for the third.

$$\left(\left| \Psi(X_u, X_u, Z_u)^T \cdot \Psi_{ABC} \right| \right)^2 = 0.25 \quad \left(\left| \Psi(X_u, X_u, Z_d)^T \cdot \Psi_{ABC} \right| \right)^2 = 0 \quad \left(\left| \Psi(X_u, X_d, Z_u)^T \cdot \Psi_{ABC} \right| \right)^2 = 0$$

$$\left(\left| \Psi(X_u, X_d, Z_d)^T \cdot \Psi_{ABC} \right| \right)^2 = 0.25 \quad \left(\left| \Psi(X_d, X_u, Z_u)^T \cdot \Psi_{ABC} \right| \right)^2 = 0 \quad \left(\left| \Psi(X_d, X_u, Z_d)^T \cdot \Psi_{ABC} \right| \right)^2 = 0.25$$

$$\left(\left| \Psi(X_d, X_d, Z_d)^T \cdot \Psi_{ABC} \right| \right)^2 = 0 \quad \left(\left| \Psi(X_d, X_d, Z_u)^T \cdot \Psi_{ABC} \right| \right)^2 = 0.25$$

$$\left(\left| \Psi(X_u, Z_u, X_u)^T \cdot \Psi_{ABC} \right| \right)^2 = 0.25 \quad \left(\left| \Psi(X_u, Z_d, X_u)^T \cdot \Psi_{ABC} \right| \right)^2 = 0 \quad \left(\left| \Psi(X_u, Z_u, X_d)^T \cdot \Psi_{ABC} \right| \right)^2 = 0$$

$$\left(\left| \Psi(X_u, Z_d, X_d)^T \cdot \Psi_{ABC} \right| \right)^2 = 0.25 \quad \left(\left| \Psi(X_d, Z_u, X_u)^T \cdot \Psi_{ABC} \right| \right)^2 = 0 \quad \left(\left| \Psi(X_d, Z_d, X_u)^T \cdot \Psi_{ABC} \right| \right)^2 = 0.25$$

$$\left(\left| \Psi(X_d, Z_d, X_d)^T \cdot \Psi_{ABC} \right| \right)^2 = 0 \quad \left(\left| \Psi(X_d, Z_u, X_d)^T \cdot \Psi_{ABC} \right| \right)^2 = 0.25$$

$$\left(\left| \Psi(Z_u, X_u, X_u)^T \cdot \Psi_{ABC} \right| \right)^2 = 0.25 \quad \left(\left| \Psi(Z_d, X_u, X_u)^T \cdot \Psi_{ABC} \right| \right)^2 = 0 \quad \left(\left| \Psi(Z_u, X_u, X_d)^T \cdot \Psi_{ABC} \right| \right)^2 = 0$$

$$\left(\left| \Psi(Z_d, X_u, X_d)^T \cdot \Psi_{ABC} \right| \right)^2 = 0.25 \quad \left(\left| \Psi(Z_u, X_d, X_u)^T \cdot \Psi_{ABC} \right| \right)^2 = 0 \quad \left(\left| \Psi(Z_d, X_d, X_u)^T \cdot \Psi_{ABC} \right| \right)^2 = 0.25$$

$$\left(\left| \Psi(Z_d, X_d, X_d)^T \cdot \Psi_{ABC} \right| \right)^2 = 0 \quad \left(\left| \Psi(Z_u, X_d, X_d)^T \cdot \Psi_{ABC} \right| \right)^2 = 0.25$$

The twelve possible outcomes are listed in two rows.

$$\begin{pmatrix} X_u X_u Z_u & X_d X_d Z_u & X_u Z_u X_u & X_d Z_u X_d & Z_u X_u X_u & Z_u X_d X_d \\ X_u X_d Z_d & X_d X_u Z_d & X_u Z_d X_d & X_d Z_d X_u & Z_d X_u X_d & Z_d X_d X_u \end{pmatrix}$$

These results are summarized as follows: if two spins have the same x-spin state then the z-spin state is u, but if they are different the z-spin state is d. Again according to the EPR criterion spin in the z-direction is an element of reality.

However, the calculations can be described in another way: (a) measurements in which three pairs have the same x spin state and (b) measurements in which one pair has the same x-spin state and the remaining pairs have different x-spin states. The following table organizes the experimental results grouped in these categories. The first two rows satisfy criterion (a) and the remaining six rows satisfy criterion (b). The left column shows the implied total x-spin state, the middle three columns how it is achieved, and the right column the z-spin state consistent with the x-spin measurement results.

XuXuXu	XuXuZu	ZuXuXu	XuZuXu	ZuZuZu
XdXdXd	XdXdZu	XdZuXd	ZuXdXd	ZuZuZu
XuXuXd	XuXuZu	XuZdXd	ZdXuXd	ZdZdZu
XdXdXu	XdXdZu	XdZdXu	ZdXdXu	ZdZdZu
XuXdXu	XuZuXu	XuXdZd	ZdXdXu	ZdZuZd
XdXuXd	XdZuXd	XdXuZd	ZdXuXd	ZdZuZd
XdXuXu	ZuXuXu	XdZdXu	XdXuZd	ZuZdZd
XuXdXd	ZuXdXd	XuZdXd	XuXdZd	ZuZdZd

As this table demonstrates, according to the x-direction spin measurements the permissible z-direction states are $Z_u Z_u Z_u$, $Z_d Z_d Z_u$, $Z_d Z_u Z_d$ and $Z_u Z_d Z_d$. This is in direct contradiction to the earlier z-direction measurements which gave the result $Z_d Z_d Z_d$, $Z_u Z_u Z_d$, $Z_u Z_d Z_u$ and $Z_d Z_d Z_u$, using the same initial wavefunction Ψ_{ABC} . In other words two measurement protocols lead to the notion that z-direction spin is an element of reality according to the EPR definition, but they dramatically disagree on the actual z-direction spin values allowed.

In the February 2001 issue of the *American Journal of Physics* on page 187, John D. Norton summarizes the situation as follows:

Two principal assumptions were made in the arguments that generated this contradiction. One was that the empirical predictions of quantum theory are reliable. The other was that the EPR reality criterion, which in turn, depends on the assumptions of separability and locality. One of these assumptions must be given up. The continuing empirical success of quantum theory has led to a consensus that it is the second assumption, the EPR reality criterion, which is to be discarded.