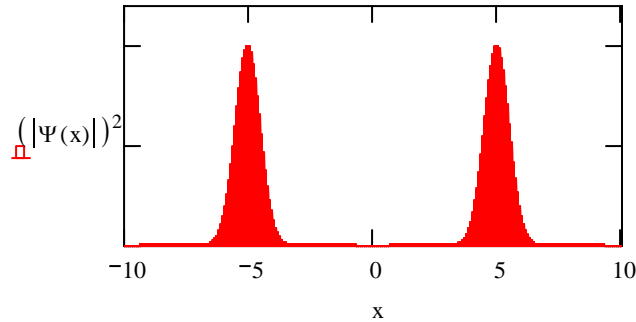


## Wigner Distribution for the Double Slit Experiment

Illumination of a double slit with a coherent incident particle beam leads to a "cat state" which can be represented by a linear superposition of two Gaussian wavepackets as shown below.

$$x := -10, -9.99 \dots 10 \quad \Psi(x) := \exp[-(x - 5)^2] + \exp[-(x + 5)^2]$$

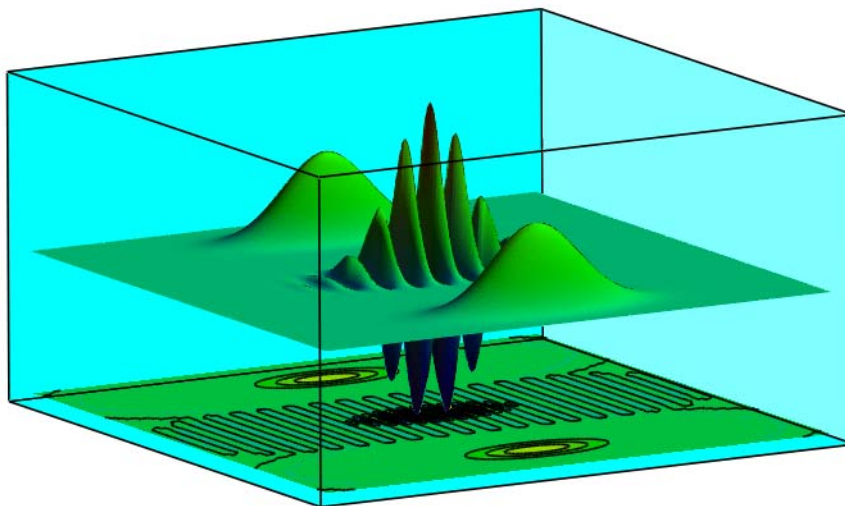


The Wigner distribution for this function is calculated and plotted below.

$$W(x,p) := \frac{1}{2\pi} \int_{-\infty}^{\infty} \left[ \exp\left[-\left(x + \frac{s}{2} - 5\right)^2\right] + \exp\left[-\left(x + \frac{s}{2} + 5\right)^2\right] \right] \cdot \exp(i \cdot p \cdot s) \cdot \left[ \exp\left[-\left(x - \frac{s}{2} - 5\right)^2\right] + \exp\left[-\left(x - \frac{s}{2} + 5\right)^2\right] \right] ds$$

$$W(x,p) := \frac{1}{\sqrt{2\pi}} \cdot \left( 2 \cdot \exp\left(-2 \cdot x^2 - \frac{1}{2} \cdot p^2\right) \cdot \cos(10 \cdot p) + \exp\left(-2 \cdot x^2 + 20 \cdot x - 50 - \frac{1}{2} \cdot p^2\right) + \exp\left(-2 \cdot x^2 - 20 \cdot x - 50 - \frac{1}{2} \cdot p^2\right) \right)$$

$$N := 100 \quad i := 0 \dots N \quad x_i := -7 + \frac{14 \cdot i}{N} \quad j := 0 \dots N \quad p_j := -6 + \frac{12 \cdot j}{N} \quad \text{Wigner}_{i,j} := W(x_i, p_j)$$



Wigner, Wigner

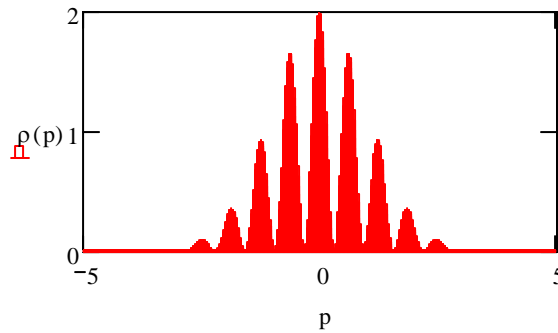
The Wigner distribution can be reconstructed from experimental measurements using quantum state tomography. Reconstructive tomography is a widely used technique in medicine, for example, for obtaining the shape of an inaccessible two-dimensional object from a set of different one-dimensional "shadows" cast by that object.

Quantum state reconstruction is possible if a system can be prepared repeatedly in the same state. Subsequent measurements on such a system are then effectively multiple measurements on the same quantum state. The theoretical Wigner distribution shown above for the double-slit experiment has been recently reconstructed for the helium atom. [See, *Nature*, **386**, 150 (1997).

Integration of the Wigner distribution for the linear superposition over the spatial coordinate yields the momentum distribution function.

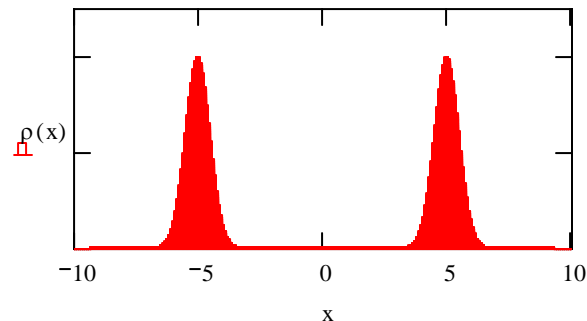
$$\rho(p) := \int_{-\infty}^{\infty} \frac{1}{\sqrt{2\pi}} \cdot \left( 2 \cdot \exp\left(-2 \cdot x^2 - \frac{1}{2} \cdot p^2\right) \cdot \cos(10 \cdot p) + \exp\left(-2 \cdot x^2 + 20 \cdot x - 50 - \frac{1}{2} \cdot p^2\right) \dots \right) dx$$

$$\rho(p) := 2 \cdot \left( \frac{1}{2} \cdot \cos(10 \cdot p) + \frac{1}{2} \right) \cdot \exp\left(\frac{-1}{2} \cdot p^2\right) \quad p := -5, -4.99 .. 5$$



Integration of the Wigner distribution for the linear superposition over the momentum coordinate returns the spatial distribution function.

$$\rho(x) := \int_{-4}^4 \frac{1}{\sqrt{2\pi}} \cdot \left( 2 \cdot \exp\left(-2 \cdot x^2 - \frac{1}{2} \cdot p^2\right) \cdot \cos(10 \cdot p) + \exp\left(-2 \cdot x^2 + 20 \cdot x - 50 - \frac{1}{2} \cdot p^2\right) \dots \right) dp \quad x := -10, -9.99 .. 10$$



See Figure 1 of "Shadows and Mirrors:Reconstructing Quantum States of Atom Motion," Physics Today, April 1998, by Leibfried, Pfau, and Monroe.

Reference: Decoherence and the Transition form Quantum to Classical, Wojciech Jurek, Physics Today, October 1991, pages 36-44.